

BRIEF COMMUNICATION

AN INVESTIGATION OF TWO-PHASE FLOW MEASUREMENT WITH ORIFICES FOR LOW-QUALITY MIXTURES

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1. INTRODUCTION

Measurement of gas-liquid two-phase flow rates is of interest in many fields of engineering, such as chemical, geothermal engineering, petroleum, power and nuclear energy. Although many methods, including ultrasonic waves, laser techniques, radiation etc., have been studied (Hewitt 1978; Reimann 1982), measurement of two-phase flow rates via orifices has received increasing attention in the last three decades. Numerous orifice equations for gas-liquid mixtures have been developed and some typical equations were proposed by Murdock (1962), James (1965), Chisholm (1967, 1974), Smith & Leang (1975) and Lin (1982). These equations can mostly be described by the following equation:

$$G = \frac{CYAFa}{\sqrt{1-\beta^4}} K_G \sqrt{2 \Delta P_{TP} \rho_G}, \quad [1]$$

where G is the mass flow rate of the gas-liquid mixture, C is the orifice discharge coefficient, Y is the compressibility coefficient of the fluid, A is the orifice flow area, Fa is the orifice thermal expansion factor, β is the ratio of the orifice diameter to the internal pipe diameter, ΔP_{TP} is the pressure drop across the orifice for the two-phase mixture, ρ_G is the density of the gas phase and K_G is called the gas phase modified coefficient, which is dependent mainly on quality and the gas-liquid density ratio, and has several different forms (proposed by different authors). Murdock (1962) employed a separated flow model and derived the following two-phase correlation equation:

$$K_G = \frac{1}{\chi + 1.26(1-\chi) \frac{C_G Y_G}{C_L} \sqrt{\frac{\rho_G}{\rho_L}}}, \quad [2]$$

where C_G and C_L are the orifice discharge coefficients for gas and liquid, respectively, Y_G is the compressibility coefficient of the gas phase and χ is the quality.

James (1965) used an effective mixture density to modify a homogeneous model, resulting in the expression

$$K_G = \sqrt{\frac{1}{\chi^{1.5} \left(1 - \frac{\rho_G}{\rho_L}\right) + \frac{\rho_G}{\rho_L}}}. \quad [3]$$

Other authors proposed different modified two-phase correlation equations based on either the separated flow model or the homogeneous flow model. But in previous works, almost all orifice equations were derived from experiments with quality $\chi > 1\%$. Some equations cannot be used particularly well and others become invalid at low quality—see Murdock's (1962) equation [2]. But low quality is widespread if the ratio of gas density to liquid density is small, even though the void fraction ϵ is large. For this reason, our work is aimed at the measurement of two-phase flow with sharp-edged orifices for low quality. The experiments were made with an air-water flow system in the quality range 0.007–1%.

2. THEORY

For low-quality gas-liquid two-phase flow, the orifice equation for metering the mass flow rate has the following form (Zhang 1985):

$$G = \frac{C_{TP} Y_{TP} A F a}{\sqrt{1 - \beta^4}} K_L \sqrt{2 \Delta P_{TP} \rho_L}, \quad [4]$$

where Y_{TP} is the compressibility coefficient of the two-phase mixture and K_L is the liquid modified coefficient. Note that in the above equation the liquid orifice discharge coefficient C_L can be used to replace the two-phase mixture orifice discharge coefficient C_{TP} if the Reynolds number for the liquid phase Re_L is greater than the critical Reynolds number Re_k (Matter *et al.* 1979). But the compressibility coefficient Y_{TP} will be different from the gas and liquid compressibility coefficients Y_G and Y_L for most cases. A simple theoretical equation for calculating Y_{TP} , proposed by Zhang (1985), is as follows:

$$Y_{TP} = Y_L(1 - \epsilon) + Y_G \epsilon, \quad [5]$$

where ϵ is the void fraction. Because $Y_L = 1$, [5] becomes

$$Y_{TP} = 1 - \epsilon + \epsilon Y_G. \quad [6]$$

It is rather difficult to obtain a simple relationship for the modified coefficient K_L from a theoretical analysis, so K_L will be determined experimentally. It is thought that K_L is a function of the density ratio ρ_G/ρ_L and of void fraction or quality, but for low quality it is better to use the void fraction as a parameter in the modified coefficient K_L instead of normal quality, because in such a range of quality the void fraction will increase or decrease significantly with a small change in quality.

3. EXPERIMENTAL APPARATUS AND RESULTS

Experiments were carried out in a two-phase flow measurement apparatus, as shown in figure 1. Air and water were used as the gas and liquid phases, respectively. Before they were mixed, air and water flow rates were measured individually. The test gas mass velocity ranged from 0.9 to 9.1 kg/m²·s with the liquid mass velocity from 840 to 1648 kg/m²·s, while the quality ranged from 0.007 to 1%. The sharp-edged orifices used for the tests were mounted on a horizontal pipe with i.d. = 50.8 mm. The diameters of the orifices were 25.36 and 21.42 mm, with a diameter ratio of $\beta = 0.499$ and $\beta = 0.422$, respectively. The pressure taps on each orifice were standard corner taps with carrier rings. Before the tests the orifices were carefully calibrated with single-phase water. During the tests the total mass flow rates and quality were obtained from the readings of the gas and liquid flowmeters. Void fraction was measured by two quick-closing valves which were installed

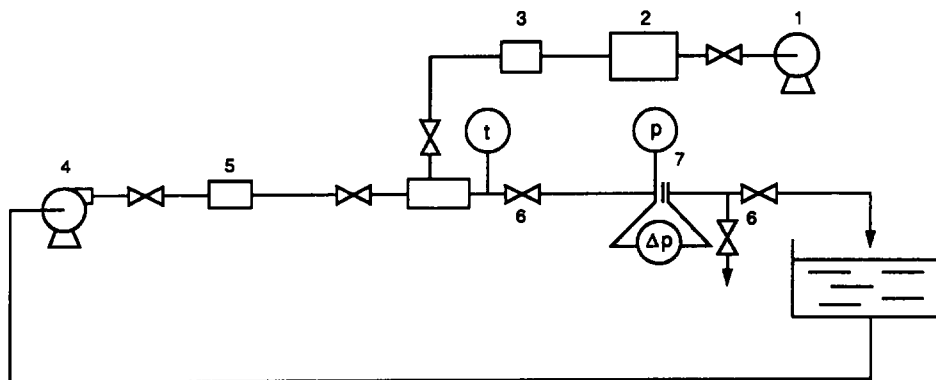


Figure 1. Schematic diagram of the experimental apparatus: 1, compressor; 2, filter; 3, gas flowmeter; 4, pump; 5, liquid flowmeter; 6, quick-closing valves; 7, orifice

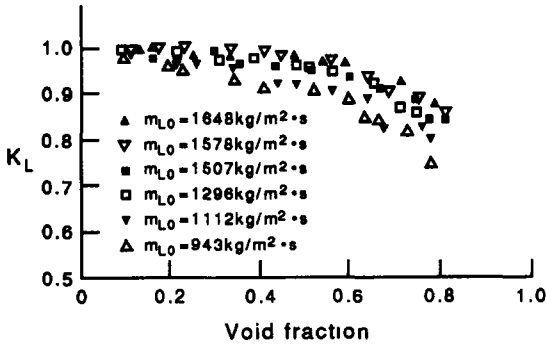


Figure 2 Experimental values of K_L vs void fraction for $\beta = 0.499$

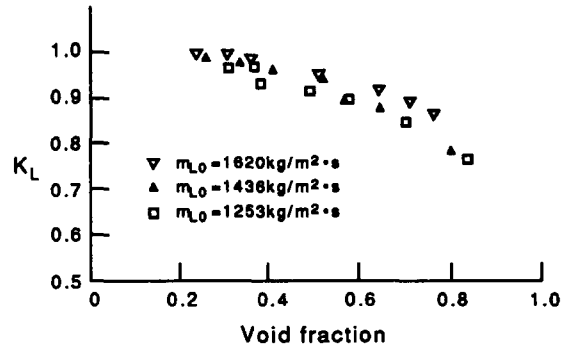


Figure 3. Experimental values of K_L vs void fraction for $\beta = 0.422$.

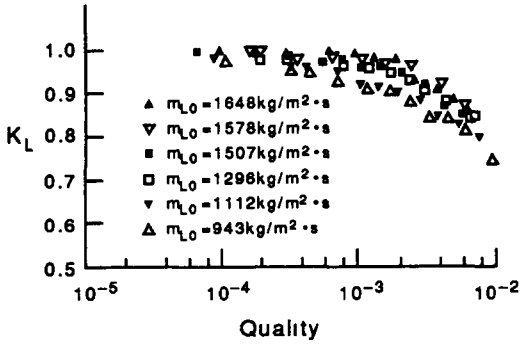


Figure 4 Experimental values of K_L vs quality for $\beta = 0.499$

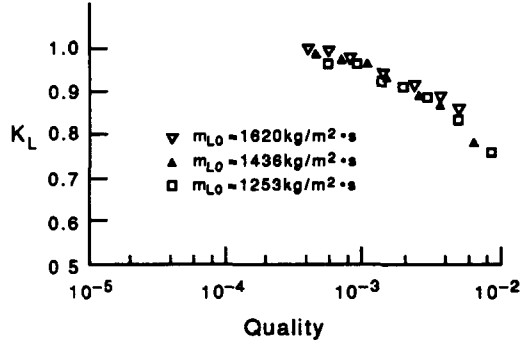


Figure 5 Experimental values of K_L vs quality for $\beta = 0.422$.

on the two sides of the sharp-edged orifice. A pressure gauge was installed before the orifice to indicate the pressure of the two-phase flow. If ΔP_{TP} is measured by a differential pressure transmitter connected to the orifice, K_L can be calculated from [4]. Since the tests were performed at normal temperature, the orifice thermal expansion factor Fa is unity and the water density is about 998 kg/m^3 . All the experimental data are summarized in table 1. Figures 2 and 3 show the experimental values of K_L vs void fraction, and figures 4 and 5 are the experimental values of K_L vs quality. In figures 2–5, m_{L0} is the liquid mass flow rate per unit area when the gas flow rate is zero; with an increase in the gas flow rate the liquid flow rate will decrease, because the position of the valve for controlling the liquid phase flow rate has not changed.

4. COMPARISON WITH RESULTS FROM PREVIOUS WORK

Because Murdock's (1962) correlation equation for two-phase flow does not apply at low quality, comparisons can only be made with the homogeneous flow model and the James model here. For the homogeneous flow model, the theoretical equation for the liquid phase modified coefficient K_L is as follows:

$$K_L = \sqrt{\frac{1}{\chi^n \left(\frac{\rho_L}{\rho_G} - 1 \right) + 1}}, \quad [7]$$

where $n = 1$. Substituting experimental data for χ , ρ_G and ρ_L into [7], the theoretical K_L can be obtained. The errors between the actual K_L and the K_L value calculated from [7] are shown in figure 6(a). Figure 6(a) shows that the actual values are greater than the calculated values and the errors increase with increasing quality.

For the James (1965) model, the equation for K_L is the same as [7], but $n = 1.5$. The deviation of actual K_L values from those given by [7] is shown in figure 6(b). It is seen from the figure that most data are negative, which means that the actual K_L is less than the calculated K_L . It is also found that the absolute errors are less than those for the homogeneous model.

Table 1. Test results

| No. | d (mm) | D (mm) | β | P_1 (MPa) | t (°C) | ρ_G (kg/m ³) | G_L (kg/s) | G_G (kg/s) | χ (%) | ϵ | ΔP_{TP} (kPa) | Y_G | Y_{TP} | α_L | K_L | G (kg/s) |
|-----|-------------|-------------|---------|----------------|-------------|----------------------------------|-----------------|-----------------|------------|------------|--------------------------|-------|----------|------------|-------|---------------|
| 1 | 25 | 36 | 50.8 | 0.499 | 0.231 | 2.44 | 3.2361 | 0.0003 | 0.0095 | 0.104 | 52.84 | 0.925 | 0.992 | 0.631 | 0.997 | 3.2364 |
| 2 | | | | | 0.235 | 2.49 | 3.2250 | 0.0006 | 0.0186 | 0.170 | 53.48 | 0.925 | 0.987 | 0.631 | 0.993 | 3.2256 |
| 3 | | | | | 0.238 | 2.52 | 3.2222 | 0.0010 | 0.0310 | 0.231 | 53.48 | 0.925 | 0.983 | 0.631 | 0.997 | 3.2232 |
| 4 | | | | | 0.242 | 2.56 | 3.2056 | 0.0020 | 0.0624 | 0.336 | 53.84 | 0.926 | 0.975 | 0.631 | 0.997 | 3.2076 |
| 5 | | | | | 0.243 | 2.57 | 3.1778 | 0.0030 | 0.0943 | 0.405 | 54.06 | 0.926 | 0.970 | 0.631 | 0.991 | 3.1808 |
| 6 | | | | | 0.247 | 2.62 | 3.1444 | 0.0041 | 0.1302 | 0.441 | 54.28 | 0.926 | 0.967 | 0.631 | 0.982 | 3.1468 |
| 7 | | | | | 0.250 | 2.65 | 3.1222 | 0.0059 | 0.1886 | 0.551 | 55.75 | 0.925 | 0.959 | 0.631 | 0.971 | 3.1281 |
| 8 | | | | | 0.255 | 2.70 | 2.9861 | 0.0079 | 0.2639 | 0.637 | 55.54 | 0.927 | 0.953 | 0.631 | 0.937 | 2.9940 |
| 9 | | | | | 0.265 | 2.82 | 2.8833 | 0.0108 | 0.3732 | 0.681 | 55.60 | 0.928 | 0.951 | 0.631 | 0.907 | 2.8971 |
| 10 | | | | | 0.275 | 2.93 | 2.7972 | 0.0143 | 0.5086 | 0.749 | 54.87 | 0.930 | 0.948 | 0.631 | 0.890 | 2.8115 |
| 11 | | | | | 0.292 | 3.12 | 2.7500 | 0.0163 | 0.5892 | 0.801 | 57.45 | 0.932 | 0.943 | 0.631 | 0.860 | 2.7663 |
| 12 | | | | | 0.231 | 2.45 | 3.1139 | 0.0005 | 0.0161 | 0.126 | 48.84 | 0.930 | 0.991 | 0.631 | 0.999 | 3.1144 |
| 13 | | | | | 0.228 | 2.42 | 3.1028 | 0.0006 | 0.0193 | 0.154 | 48.62 | 0.930 | 0.989 | 0.631 | 1.000 | 3.1034 |
| 14 | | | | | 0.230 | 2.44 | 3.0250 | 0.0011 | 0.0363 | 0.248 | 48.84 | 0.930 | 0.983 | 0.631 | 0.983 | 3.0261 |
| 15 | | | | | 0.232 | 2.46 | 3.0194 | 0.0020 | 0.0662 | 0.330 | 48.47 | 0.930 | 0.977 | 0.631 | 0.987 | 3.0214 |
| 16 | | | | | 0.237 | 2.52 | 2.9778 | 0.0032 | 0.1073 | 0.476 | 48.40 | 0.932 | 0.968 | 0.631 | 0.984 | 2.9810 |
| 17 | | | | | 0.240 | 2.55 | 2.9361 | 0.0045 | 0.1530 | 0.536 | 48.84 | 0.932 | 0.964 | 0.631 | 0.971 | 2.9406 |
| 18 | | | | | 0.245 | 2.61 | 2.9194 | 0.0071 | 0.2426 | 0.586 | 48.92 | 0.933 | 0.961 | 0.631 | 0.968 | 2.9425 |
| 19 | | | | | 0.257 | 2.74 | 2.8250 | 0.0113 | 0.3984 | 0.705 | 50.75 | 0.933 | 0.953 | 0.631 | 0.929 | 2.8363 |
| 20 | | | | | 0.273 | 2.91 | 2.6972 | 0.0169 | 0.5860 | 0.779 | 52.15 | 0.935 | 0.949 | 0.631 | 0.880 | 2.7131 |
| 21 | | | | | 0.216 | 2.26 | 2.9694 | 0.0002 | 0.0067 | 0.086 | 44.36 | 0.932 | 0.994 | 0.631 | 0.997 | 2.9696 |
| 22 | | | | | 0.223 | 2.39 | 2.9222 | 0.0009 | 0.0308 | 0.204 | 44.13 | 0.933 | 0.986 | 0.631 | 0.992 | 2.9231 |
| 23 | | | | | 0.226 | 2.42 | 2.8972 | 0.0016 | 0.0552 | 0.305 | 45.23 | 0.933 | 0.980 | 0.631 | 0.978 | 2.8988 |
| 24 | | | | | 0.228 | 2.44 | 2.8722 | 0.0021 | 0.0731 | 0.382 | 44.50 | 0.933 | 0.974 | 0.631 | 0.965 | 2.8743 |
| 25 | | | | | 0.230 | 2.47 | 2.8361 | 0.0030 | 0.1057 | 0.467 | 45.53 | 0.933 | 0.969 | 0.631 | 0.965 | 2.8391 |
| 26 | | | | | 0.231 | 2.48 | 2.8167 | 0.0041 | 0.1453 | 0.501 | 45.53 | 0.933 | 0.966 | 0.631 | 0.961 | 2.8208 |
| 27 | | | | | 0.239 | 2.57 | 2.7944 | 0.0057 | 0.2036 | 0.559 | 45.97 | 0.935 | 0.964 | 0.631 | 0.953 | 2.8001 |
| 28 | | | | | 0.245 | 2.63 | 2.7361 | 0.0082 | 0.2988 | 0.649 | 47.44 | 0.934 | 0.957 | 0.631 | 0.975 | 2.7443 |
| 29 | | | | | 0.250 | 2.68 | 2.6278 | 0.0107 | 0.4055 | 0.707 | 49.65 | 0.933 | 0.953 | 0.631 | 0.872 | 2.6355 |
| 30 | | | | | 0.255 | 2.73 | 2.5944 | 0.0148 | 0.5672 | 0.746 | 50.38 | 0.933 | 0.950 | 0.632 | 0.859 | 2.6092 |
| 31 | | | | | 0.203 | 2.15 | 2.5972 | 0.0005 | 0.0192 | 0.155 | 35.16 | 0.941 | 0.991 | 0.631 | 0.983 | 2.5977 |
| 32 | | | | | 0.208 | 2.25 | 2.5972 | 0.0008 | 0.0308 | 0.206 | 35.38 | 0.941 | 0.988 | 0.631 | 0.983 | 2.5980 |
| 33 | | | | | 0.211 | 2.28 | 2.5833 | 0.0014 | 0.0542 | 0.296 | 35.68 | 0.941 | 0.983 | 0.632 | 0.977 | 2.5947 |
| 34 | | | | | 0.216 | 2.33 | 2.5444 | 0.0020 | 0.0785 | 0.347 | 35.68 | 0.943 | 0.980 | 0.632 | 0.965 | 2.5464 |
| 35 | | | | | 0.218 | 2.35 | 2.5444 | 0.0030 | 0.1182 | 0.429 | 36.12 | 0.943 | 0.976 | 0.632 | 0.961 | 2.5391 |
| 36 | | | | | 0.221 | 2.39 | 2.5025 | 0.0043 | 0.1715 | 0.507 | 36.41 | 0.943 | 0.971 | 0.632 | 0.950 | 2.5071 |
| 37 | | | | | 0.228 | 2.47 | 2.4667 | 0.0059 | 0.2386 | 0.597 | 36.63 | 0.944 | 0.967 | 0.632 | 0.938 | 2.4726 |
| 38 | | | | | 0.233 | 2.53 | 2.4333 | 0.0074 | 0.3032 | 0.664 | 37.88 | 0.944 | 0.963 | 0.632 | 0.914 | 2.4407 |
| 39 | | | | | 0.242 | 2.63 | 2.3659 | 0.0104 | 0.4380 | 0.742 | 38.62 | 0.944 | 0.958 | 0.632 | 0.885 | 2.3743 |

| | | | | | | | | | | | | | |
|----|-------|----|------|--------|--------|--------|-------|--------|-------|-------|-------|-------|--------|
| 40 | 0.251 | 30 | 2.72 | 2.2944 | 0.0148 | 0.6409 | 0.766 | 40.09 | 0.944 | 0.957 | 0.632 | 0.846 | 2.3092 |
| 41 | 0.251 | 35 | 2.68 | 2.2944 | 0.0159 | 0.6882 | 0.801 | 40.01 | 0.944 | 0.955 | 0.632 | 0.849 | 2.3103 |
| 42 | 0.185 | 35 | 1.99 | 2.2528 | 0.0089 | 0.0089 | 0.110 | 35.62 | 0.953 | 0.995 | 0.633 | 0.988 | 2.2530 |
| 43 | 0.194 | 30 | 2.13 | 2.1722 | 0.0007 | 0.0322 | 0.212 | 25.37 | 0.955 | 0.990 | 0.633 | 0.965 | 2.2530 |
| 44 | 0.194 | 30 | 2.13 | 2.1750 | 0.0009 | 0.0414 | 0.257 | 25.45 | 0.955 | 0.988 | 0.633 | 0.967 | 2.1759 |
| 45 | 0.201 | 30 | 2.21 | 2.1389 | 0.0015 | 0.0701 | 0.337 | 25.38 | 0.957 | 0.986 | 0.633 | 0.955 | 2.1404 |
| 46 | 0.203 | 30 | 2.22 | 2.1194 | 0.0022 | 0.1037 | 0.433 | 27.00 | 0.955 | 0.981 | 0.633 | 0.923 | 2.1216 |
| 47 | 0.206 | 30 | 2.26 | 2.1278 | 0.0030 | 0.1408 | 0.472 | 27.51 | 0.955 | 0.979 | 0.633 | 0.920 | 2.1308 |
| 48 | 0.206 | 30 | 2.25 | 2.1028 | 0.0040 | 0.1899 | 0.559 | 27.80 | 0.954 | 0.974 | 0.633 | 0.909 | 2.1068 |
| 49 | 0.211 | 30 | 2.31 | 2.0750 | 0.0059 | 0.2835 | 0.634 | 28.63 | 0.954 | 0.971 | 0.633 | 0.888 | 2.0809 |
| 50 | 0.216 | 30 | 2.35 | 2.0333 | 0.0077 | 0.3773 | 0.669 | 30.28 | 0.951 | 0.971 | 0.633 | 0.850 | 2.0410 |
| 51 | 0.228 | 30 | 2.49 | 1.9917 | 0.0109 | 0.5443 | 0.753 | 30.30 | 0.955 | 0.966 | 0.633 | 0.836 | 2.0026 |
| 52 | 0.228 | 30 | 2.50 | 1.9194 | 0.0148 | 0.7652 | 0.768 | 30.53 | 0.955 | 0.965 | 0.634 | 0.802 | 1.9342 |
| 53 | 0.179 | 35 | 1.96 | 1.8583 | 0.0002 | 0.0108 | 0.096 | 18.02 | 0.956 | 0.997 | 0.634 | 0.972 | 1.8585 |
| 54 | 0.184 | 35 | 2.01 | 1.8306 | 0.0006 | 0.0328 | 0.192 | 18.09 | 0.966 | 0.993 | 0.634 | 0.959 | 1.8312 |
| 55 | 0.188 | 32 | 2.08 | 1.8056 | 0.0008 | 0.0433 | 0.225 | 18.02 | 0.967 | 0.993 | 0.634 | 0.948 | 1.8064 |
| 56 | 0.193 | 32 | 2.13 | 1.7944 | 0.0013 | 0.0724 | 0.343 | 18.76 | 0.967 | 0.989 | 0.634 | 0.929 | 1.7957 |
| 57 | 0.193 | 32 | 2.13 | 1.7917 | 0.0021 | 0.1171 | 0.404 | 19.50 | 0.966 | 0.986 | 0.634 | 0.911 | 1.7938 |
| 58 | 0.194 | 32 | 2.14 | 1.7750 | 0.0030 | 0.1687 | 0.520 | 19.42 | 0.966 | 0.982 | 0.634 | 0.909 | 1.7780 |
| 59 | 0.196 | 32 | 2.16 | 1.7778 | 0.0043 | 0.2413 | 0.596 | 20.74 | 0.965 | 0.979 | 0.634 | 0.884 | 1.7821 |
| 60 | 0.198 | 32 | 2.17 | 1.7361 | 0.0058 | 0.3330 | 0.632 | 21.70 | 0.963 | 0.977 | 0.635 | 0.846 | 1.7419 |
| 61 | 0.203 | 32 | 2.23 | 1.7000 | 0.0079 | 0.4626 | 0.662 | 21.26 | 0.955 | 0.970 | 0.635 | 0.843 | 1.7079 |
| 62 | 0.211 | 32 | 2.31 | 1.7278 | 0.0104 | 0.5983 | 0.726 | 23.17 | 0.963 | 0.973 | 0.635 | 0.820 | 1.7382 |
| 63 | 0.216 | 32 | 2.36 | 1.6556 | 0.0159 | 0.9512 | 0.774 | 26.11 | 0.958 | 0.967 | 0.635 | 0.747 | 1.6715 |
| 64 | 0.294 | 21 | 2.96 | 3.2472 | 0.0013 | 0.0400 | 0.240 | 111.99 | 0.878 | 0.971 | 0.623 | 0.998 | 3.2485 |
| 65 | 0.289 | 21 | 2.90 | 3.2250 | 0.0018 | 0.0558 | 0.304 | 113.36 | 0.870 | 0.960 | 0.623 | 0.996 | 3.2268 |
| 66 | 0.294 | 21 | 2.91 | 3.2639 | 0.0027 | 0.0827 | 0.357 | 122.19 | 0.863 | 0.951 | 0.623 | 0.981 | 3.2666 |
| 67 | 0.299 | 21 | 2.93 | 3.1806 | 0.0045 | 0.1413 | 0.505 | 130.62 | 0.860 | 0.929 | 0.623 | 0.946 | 3.1851 |
| 68 | 0.304 | 21 | 2.95 | 3.0806 | 0.0075 | 0.2429 | 0.641 | 137.69 | 0.856 | 0.908 | 0.623 | 0.915 | 3.0881 |
| 69 | 0.314 | 21 | 3.07 | 2.9639 | 0.0109 | 0.3664 | 0.710 | 137.69 | 0.860 | 0.901 | 0.623 | 0.888 | 2.9748 |
| 70 | 0.324 | 21 | 3.07 | 2.8639 | 0.0145 | 0.5037 | 0.759 | 137.69 | 0.860 | 0.894 | 0.623 | 0.866 | 2.8784 |
| 71 | 0.260 | 21 | 2.66 | 2.8472 | 0.0013 | 0.0456 | 0.259 | 88.26 | 0.888 | 0.971 | 0.623 | 0.985 | 2.8485 |
| 72 | 0.260 | 21 | 2.66 | 2.8056 | 0.0020 | 0.0712 | 0.332 | 89.24 | 0.887 | 0.962 | 0.623 | 0.974 | 2.8076 |
| 73 | 0.260 | 21 | 2.65 | 2.7583 | 0.0030 | 0.1086 | 0.407 | 90.42 | 0.887 | 0.954 | 0.623 | 0.961 | 2.7613 |
| 74 | 0.265 | 21 | 2.68 | 2.7222 | 0.0041 | 0.1504 | 0.519 | 96.30 | 0.882 | 0.939 | 0.623 | 0.934 | 2.7263 |
| 75 | 0.270 | 21 | 2.73 | 2.6139 | 0.0068 | 0.2595 | 0.569 | 97.87 | 0.882 | 0.933 | 0.623 | 0.896 | 2.6207 |
| 76 | 0.278 | 21 | 2.81 | 2.5472 | 0.0093 | 0.3638 | 0.641 | 100.62 | 0.882 | 0.924 | 0.623 | 0.870 | 2.5565 |
| 77 | 0.287 | 21 | 2.95 | 2.3500 | 0.0152 | 0.6435 | 0.799 | 109.25 | 0.879 | 0.908 | 0.623 | 0.866 | 2.3652 |
| 78 | 0.240 | 21 | 2.51 | 2.5000 | 0.0014 | 0.0560 | 0.310 | 71.20 | 0.902 | 0.970 | 0.623 | 0.965 | 2.5014 |
| 79 | 0.245 | 21 | 2.57 | 2.5000 | 0.0023 | 0.0919 | 0.363 | 71.98 | 0.902 | 0.964 | 0.623 | 0.965 | 2.5023 |
| 80 | 0.245 | 21 | 2.56 | 2.4583 | 0.0034 | 0.1381 | 0.381 | 76.49 | 0.899 | 0.962 | 0.623 | 0.924 | 2.4617 |
| 81 | 0.250 | 21 | 2.59 | 2.4528 | 0.0049 | 0.1994 | 0.490 | 79.24 | 0.900 | 0.951 | 0.623 | 0.916 | 2.4577 |
| 82 | 0.260 | 21 | 2.70 | 2.3972 | 0.0074 | 0.3077 | 0.584 | 81.59 | 0.899 | 0.941 | 0.623 | 0.893 | 2.4046 |
| 83 | 0.270 | 21 | 2.79 | 2.3056 | 0.0113 | 0.4877 | 0.700 | 86.89 | 0.897 | 0.928 | 0.623 | 0.845 | 2.3169 |
| 84 | 0.279 | 21 | 2.88 | 2.1000 | 0.0179 | 0.8452 | 0.837 | 91.40 | 0.895 | 0.915 | 0.623 | 0.764 | 2.1179 |

21 42 50 8 0 422

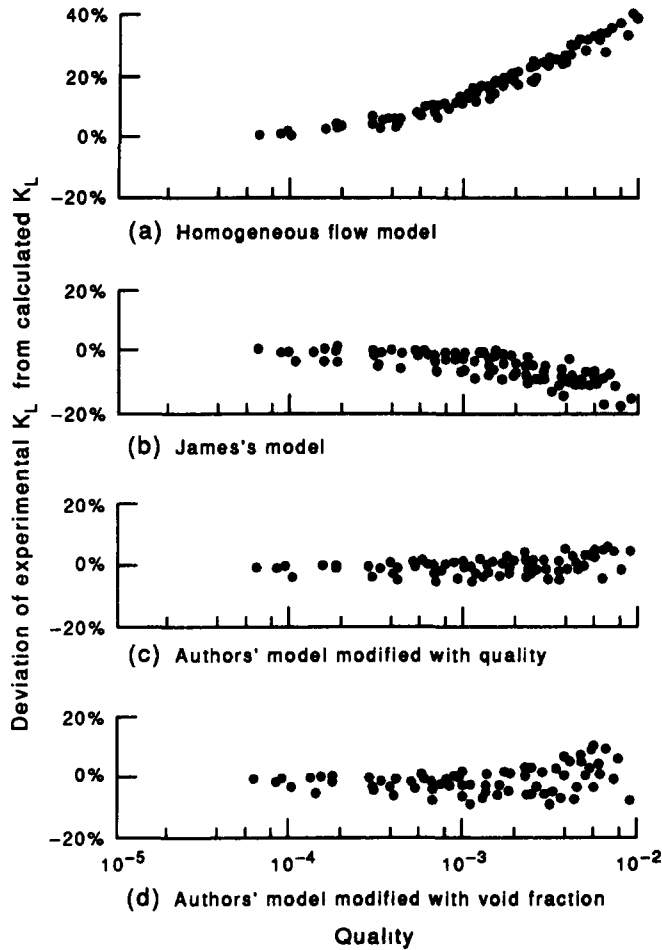


Figure 6 Deviation of experimental K_L values from those from different models.

5. NEW EQUATIONS

Because most errors between actual and calculated K_L values are greater than zero for the homogeneous flow model, but less than zero for the James (1965) model, it is thought that [7] may be appropriate for two-phase flow with n being from 1 to 1.5. Using the experimental data shown in table 1, n is found to be

$$n = 1.25 + 0.25\sqrt[3]{\chi}. \quad [8]$$

When χ approaches unity, $n = 1.5$, i.e. it becomes the same as the James model. If [8] is used to calculate K_L , we find it is quite close to the experimental K_L shown in figure 6(c).

If the void fraction is employed as a parameter in the modified coefficient K_L then, from the homogeneous flow model, the theoretical equation is

$$K_L = \sqrt{\epsilon^{n_\epsilon} \left(\frac{\rho_G}{\rho_L} - 1 \right) + 1}, \quad [9]$$

where $n_\epsilon = 1$. But this equation has significant errors compared with the experimental data. To fit the data, the method of least squares was used and we obtained $n_\epsilon = 4$. Figure 6(d) shows the differences between the actual K_L values and those calculated from [9].

6. DISCUSSION AND CONCLUSIONS

For a gas-liquid mixture of low quality, the two-phase flow rate can be measured with sharp-edged orifices. If the quality is used to modify the orifice equation, then the total mass flow rate of the two-phase mixture can be calculated using the following equation:

$$G = \frac{C_L Y_{TP} A F a}{\sqrt{1 - \beta^4}} \sqrt{\frac{2 \Delta P_{TP} \rho_L}{\chi^n \left(\frac{\rho_L}{\rho_G} - 1 \right) + 1}}, \quad [10]$$

where $n = 1.25 + 0.25 \sqrt[3]{\chi}$.

The orifice equation for the single-phase flow rate may be modified using the void fraction when the quality of the mixture is low, because under these conditions the void fraction will change more sharply than the quality and have greater values. Introducing void fraction, the modified orifice equation becomes

$$G = \frac{C_L Y_{TP} A F a}{\sqrt{1 - \beta^4}} \sqrt{2 \Delta P_{TP} \rho_L \left[\epsilon^4 \left(\frac{\rho_G}{\rho_L} - 1 \right) + 1 \right]}. \quad [11]$$

Although the deviation of experimental K_L values from those calculated by [9] is not sufficiently small [see figure 6(d)], [9] is very useful for low quality. If the void fraction is estimated, the liquid modified coefficient may be evaluated quickly. Moreover, for $\epsilon \leq 0.5$, then from [9] we get $K_L \leq 0.96$, which means that there is no significant difference between the modified orifice equation for the two-phase flow rate and the common orifice equation for the single-phase liquid flow rate. Furthermore, if the compressibility coefficient Y_{TP} is considered, we can expect that the maximum error in the mass flow rate will be within 10% if the orifice equation of single-phase liquid flow is used to measure the two-phase mass flow rate, because for most cases the compressibility coefficient is < 0.95 .

The above conclusions are derived from experiments with an air-water two-phase flow system. To contain their general validity additional experiments with different two-phase mixtures, pipe diameters and density ratios are required. However, the present work may be instructive in the further study of the measurement of gas-liquid mixture at low quality.

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